

# NECESSARY OPTIMALITY CONDITIONS FOR AN OPTIMAL CONTROL PROBLEM OF 2D $g$ -NAVIER-STOKES EQUATIONS

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## Abstract

Considered here is the optimal control problem of 2D  $g$ -Navier-Stokes equations. We prove the existence of optimal solutions and establish the first-order necessary optimality conditions.

Chúng tôi nghiên cứu bài toán điều khiển tối ưu của hệ phương trình  $g$ -Navier-Stokes 2 chiều. Chúng tôi chứng minh sự tồn tại nghiệm tối ưu và thiết lập các điều kiện tối ưu cần cấp một.

## Index terms

$g$ -Navier-Stokes equations; existence of optimal solution; first-order necessary optimality conditions.

## 1. Introduction

Let  $\Omega$  be a bounded domain in  $\mathbb{R}^2$  with boundary  $\Gamma$ . In this paper we consider the following 2D  $g$ -Navier-Stokes equations:

$$\begin{cases} \frac{\partial y}{\partial t} - \nu \Delta y + (y \cdot \nabla)y + \nabla p & = f \text{ in } (0, T) \times \Omega, \\ \nabla \cdot (gy) & = 0 \text{ in } (0, T) \times \Omega, \\ y(0, x) & = y_0, x \in \Omega, \end{cases} \quad (1.1)$$

where  $y = y(x, t) = (y_1, y_2)$  is the unknown velocity vector,  $p = p(x, t)$  is the unknown pressure,  $\nu > 0$  is the kinematic viscosity coefficient,  $y_0$  is the initial velocity, the external force  $f$  plays the role of control.

The  $g$ -Navier-Stokes equations is a variation of the standard Navier-Stokes equations. More precisely, when  $g \equiv \text{const}$  we get the usual Navier-Stokes equations. The 2D  $g$ -Navier-Stokes equations arise in a natural way when we study the standard 3D problem in thin domains. We refer the reader to [6] for a derivation of the 2D  $g$ -Navier-Stokes equations from the 3D Navier-Stokes equations and a relationship between them. As mentioned in [6], good properties of the 2D  $g$ -Navier-Stokes equations can lead to an

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initiate the study of the Navier-Stokes equations on the thin three dimensional domain  $\Omega_g = \Omega \times (0, g)$ . In the last few years, the existence and asymptotic behavior of weak solutions to 2D  $g$ -Navier-Stokes equations have been studied extensively in [1], [4], [6]. In a recent work [2], the authors proved the existence and numerical approximation of strong solutions to the 2D  $g$ -Navier-Stokes equations. The long-time behavior of the strong solutions was studied more recently in [3], [5].

In this paper we consider an optimal control with a quadratic objective functional for 2D  $g$ -Navier-Stokes equations (see Section 3 for details). To do this, we assume that the function  $g$  satisfies the following assumption:

(G)  $g \in W^{1,\infty}(\Omega)$  such that:

$$0 < m_0 \leq g(x) \leq M_0 \forall x = (x_1, x_2) \in \Omega, \text{ and } |\nabla g|_\infty < m_0 \lambda_1^{1/2},$$

where  $\lambda_1 > 0$  is the first eigenvalue of the  $g$ -Stokes operator in  $\Omega$  (i.e. the operator  $A$  defined in Section 2 below).

The paper is organized as follows. In Section 2, for convenience of the reader, we recall some auxiliary results on the 2D  $g$ -Navier-Stokes equations, which will be used later. Section 3 proves the existence of optimal solutions. The first-order necessary optimality condition is proved in the last section.

## 2. Preliminary results

### 2.1. Function spaces and operators

Let  $L^2(\Omega, g) = (L^2(\Omega))^2$  and  $H_0^1(\Omega, g) = (H_0^1(\Omega))^2$  be endowed, respectively, with the inner products

$$(u, v)_g = \int_{\Omega} u \cdot v g dx, \quad u, v \in L^2(\Omega, g),$$

and

$$((u, v))_g = \int_{\Omega} \sum_{j=1}^2 \nabla u_j \cdot \nabla v_j g dx, \quad u = (u_1, u_2), v = (v_1, v_2) \in H_0^1(\Omega, g),$$

and norms  $|u|^2 = (u, u)_g$ ,  $\|u\|^2 = ((u, u))_g$ . Thanks to assumption (G), the norms  $|\cdot|$  and  $\|\cdot\|$  are equivalent to the usual ones in  $(L^2(\Omega))^2$  and in  $(H_0^1(\Omega))^2$ .

Let

$$\mathcal{V} = \{u \in (C_0^\infty(\Omega))^2 : \nabla \cdot (gu) = 0\}.$$

Denote by  $H_g$  the closure of  $\mathcal{V}$  in  $L^2(\Omega, g)$ , and by  $V_g$  the closure of  $\mathcal{V}$  in  $H_0^1(\Omega, g)$ . It follows that  $V_g \subset H_g \equiv H'_g \subset V'_g$ , where the injections are dense and continuous. We will use  $\|\cdot\|_*$  for the norm in  $V'_g$ , and  $\langle \cdot, \cdot \rangle$  for duality pairing between  $V_g$  and  $V'_g$ . We refer the reader to [7], [8] for more details.

To deal with the time derivative in the state equation, we introduce the common spaces of functions  $y$  whose time derivatives  $y_t$  exists as abstract functions

$$W^\alpha(0, T; V_g) := \{y \in L^2(0, T; V_g) : y_t \in L^\alpha(0, T; V'_g)\},$$

$$W(0, T) := W^2(0, T; V_g).$$

where  $1 \leq \alpha \leq 2$ . Endowed with the norms

$$\|y\|_{W^\alpha} := \|y\|_{W^\alpha(0, T; V_g)} = \|y\|_{L^2(V_g)} + \|y_t\|_{L^\alpha(V'_g)},$$

$$\|y\|_W := \|y\|_{W^2},$$

these spaces are Banach spaces.

Set  $A : V_g \rightarrow V'_g$  by  $\langle Au, v \rangle = ((u, v))_g$ ; Denote  $D(A) = \{u \in V_g : Au \in H_g\}$ , then  $D(A) = H^2(\Omega, g) \cap V_g$  and  $Au = -P_g \Delta u \quad \forall u \in D(A)$ , where  $P_g$  is the ortho-projector from  $L^2(\Omega, g)$  onto  $H_g$ .

Set  $B : V_g \times V_g \rightarrow V'_g$  by  $\langle B(u, v), w \rangle = b(u, v, w)$ , where

$$b(u, v, w) = \sum_{j,k=1}^2 \int_{\Omega} u_j \frac{\partial v_k}{\partial x_j} w_k g dx,$$

whenever the integrals make sense. It is easy to check that if  $u, v, w \in V_g$ , then

$$b(u, v, w) = -b(u, w, v).$$

Hence

$$b(u, v, v) = 0 \quad \forall u, v \in V_g.$$

Let  $u \in L^2(0, T; V_g)$ , then the function  $Cu$  defined by

$$(Cu(t), v)_g = ((\frac{\nabla g}{g} \cdot \nabla)u, v)_g = b(\frac{\nabla g}{g}, u, v), \forall v \in V_g,$$

belongs to  $L^2(0, T; H_g)$ , and hence also belongs to  $L^2(0, T; V'_g)$ . Since

$$-\frac{1}{g}(\nabla \cdot g \nabla)u = -\Delta u - (\frac{\nabla g}{g} \cdot \nabla)u,$$

we have

$$(-\Delta u, v)_g = ((u, v))_g + ((\frac{\nabla g}{g} \cdot \nabla)u, v)_g = (Au, v)_g + ((\frac{\nabla g}{g} \cdot \nabla)u, v)_g, \quad \forall u, v \in V_g.$$

**Definition 2.1** (Weak solution). *Let  $f \in L^2(0, T; V'_g)$  and  $y_0 \in H_g$  be given. A function  $y \in L^2(0, T; V_g)$  with  $y_t \in L^2(0, T; V'_g)$ , i.e.  $y \in W(0, T)$ , is called weak solution of problem (1.1) if it fulfills*

$$y_t + \nu Ay + \nu Cy + B(y) = f \quad \text{in } L^2(0, T; V'_g),$$

$$y(0) = y_0 \quad \text{in } H_g. \tag{2.1}$$

The following result was proved in [1].

**Theorem 2.1** (Existence and uniqueness of solutions). *Let  $\Omega$  be a bounded and locally Lipschitz domain in  $\mathbb{R}^2$ . Then for every  $f \in L^2(0, T; V'_g)$  and  $y_0 \in H_g$ , the equation (2.1) has a unique weak solution  $y \in W(0, T)$ .*

## 2.2. Linearized equations

We will need some results about the linearized equations. Given a state  $\bar{y} \in W(0, T)$ , we consider the system

$$\begin{aligned} y_t + \nu Ay + \nu Cy + B'(\bar{y})y &= f \quad \text{in } L^2(0, T; V'_g), \\ y(0) &= y_0 \quad \text{in } H_g. \end{aligned} \quad (2.2)$$

Here,  $B'(\bar{y})y$  denotes the Fréchet derivative of  $B$  with respect to the state  $\bar{y}$ . It is itself a functional of  $L^2(0, T; V'_g)$ , which for  $v \in L^2(0, T; V_g)$  is given by

$$(B'(\bar{y})y, v) = \int_0^T (b(\bar{y}(t), y(t), v(t)) + b(y(t), \bar{y}(t), v(t))) dt. \quad (2.3)$$

**Lemma 2.1.** [11] *The operator  $B : W(0, T) \rightarrow L^2(0, T; V'_g)$  is twice Fréchet differentiable. All derivatives of third or higher order vanish. The first derivative is given by (2.3). It can be estimated as*

$$\|B'(\bar{y})y\|_{L^2(V'_g)} \leq c\|\bar{y}\|_W\|y\|_W.$$

As for quadratic functions, the second derivative is independent of  $\bar{y}$ :

$$(B''(\bar{y})[y_1, y_2], v) = \int_0^T (b(y_1(t), y_2(t), v(t)) + b(y_2(t), y_1(t), v(t))) dt.$$

The adjoint of  $B'(\bar{y})$ , called  $B'(\bar{y})^*$ , is a linear and continuous operator from  $L^2(0, T; V_g)$  to  $W^*(0, T)$ . It can be written as

$$(B'(\bar{y})^*\lambda, w) = \int_0^T (b(\bar{y}(t), w(t), \lambda(t)) + b(w(t), \bar{y}(t), \lambda(t))) dt.$$

**Theorem 2.2.** [11] *Let  $f \in L^2(0, T; V'_g)$ ,  $y_0 \in H_g$  and  $\bar{y} \in W(0, T)$  be given. Then the equation (2.2) has a unique weak solution  $y \in W(0, T)$ .*

## 2.3. The control-to-state mapping

We will study the mapping: right-hand side  $\mapsto$  solution, the so-called *control-to-state mapping*. Let  $u \in L^2(Q)^2$  denote the control, then we will use  $u$  as  $f$  in (1.1).

**Definition 2.2** (Solution mapping). *Consider the system (1.1). The mapping  $u \mapsto y$ , where  $y$  is the weak solution of (1.1) with the control right-hand side  $u$  and fixed initial value  $y_0$ , is denoted by  $S$ , i.e.  $y = S(u)$ .*

**Lemma 2.2.** [11] *The control-to-state mapping is Fréchet differentiable as mapping from  $L^2(0, T; V'_g)$  to  $W(0, T)$ . The derivative at  $\bar{u} \in L^2(0, T; V'_g)$  in direction  $h \in L^2(0, T; V'_g)$  is given by  $S'(\bar{u})h = y$ , where  $y$  is the weak solution of*

$$\begin{aligned} y_t + \nu Ay + \nu Cy + B'(\bar{y})y &= h \quad \text{in } L^2(0, T; V'_g), \\ y(0) &= 0 \quad \text{in } H_g, \end{aligned}$$

with  $\bar{y} = S(\bar{u})$ .

In order to establish first-order optimality conditions, we need the adjoint operator of  $S'(\bar{u})$  denote by  $S'(\bar{u})^*$ . By Lemma 2.2, we can regard  $S'(\bar{u})$  as linear operator from  $L^2(0, T; V'_g)$  to  $W_0$ , where  $W_0$  is defined as a closed linear subspace of  $W(0, T)$  by

$$W_0 := \{y \in W(0, T) : y(0) = 0\}.$$

Hence, the adjoint will be a mapping from  $W_0^*$  to  $L^2(0, T; V_g)$ .

**Lemma 2.3.** [11] *Let be  $\bar{u} \in L^2(Q)^2$ . Then the operator  $S'(\bar{u})^*$  is linear and continuous from  $W_0^*$  to  $L^2(0, T; V_g)$ . Its action is defined as follows. Take  $z$  in  $W_0^*$ , then  $\lambda = S'(\bar{u})^*z$  holds if and only if*

$$\langle w_t + \nu Aw + \nu Cw + B'(\bar{y})w, \lambda \rangle_{L^2(V'_g), L^2(V_g)} = \langle z, w \rangle_{W_0^*, W_0} \quad (2.4)$$

for all  $w \in W_0$ .

**Lemma 2.4.** [11] *Let be  $\bar{u} \in L^2(Q)^2$  given. Suppose the right-hand side  $z$  of (2.4) is in the form  $z = z_1 + z_2$  with functionals  $z_1 \in L^{4/3}(0, T; V'_g) \cap W_0^*$  and  $z_2 \in W_0^*$  defined by  $z_2(w) = (z_T, w(T))$ ,  $z_T \in H_g$ . Then  $\lambda = S'(\bar{u})^*z$  is the weak solution of*

$$\begin{aligned} -\lambda_t + \nu A\lambda + \nu C\lambda + B'(\bar{y})^*\lambda &= z_1 \quad \text{in } L^{4/3}(0, T; V'_g), \\ \lambda(T) &= z_T. \end{aligned}$$

Furthermore, it holds  $\lambda \in W^{4/3}(0, T)$ .

### 3. The optimal control problem

We are considering optimal control of the instationary  $g$ -Navier-Stokes equations. The minimization of the following quadratic objective functional serves as model problem:

$$J(y, u) = \frac{\alpha_T}{2} \int_{\Omega} |y(x, T) - y_T(x)|^2 dx + \frac{\gamma}{2} \int_Q |u(x, t)|^2 dx dt. \quad (3.1)$$

They are weighted by the coefficients  $\alpha_T, \gamma$ . The free variables-state  $y$  and control  $u$  have to fulfill the instationary 2D  $g$ -Navier-Stokes equations

$$\begin{cases} y_t - \nu \Delta y + (y \cdot \nabla)y + \nabla p &= u \text{ in } Q, \\ \nabla \cdot (gy) &= 0 \text{ in } Q, \\ y &= 0 \text{ on } \Sigma, \\ y(0, x) &= y_0 \text{ in } \Omega. \end{cases} \quad (3.2)$$

The control has to satisfy inequality constraints:

$$u_{a,i}(x, t) \leq u_i(x, t) \leq u_{b,i}(x, t) \quad \text{a.e. on } Q, \quad i = 1, 2.$$

### 3.1. Setting of the problem

Let us specify the problem setting. Unless other conditions are imposed, we assume that the ingredients of the optimal problem satisfy the following:

- (i) The domain  $\Omega$  is supposed to be an open bounded subset of  $\mathbb{R}^2$  with Lipschitz boundary  $\Gamma$ . We denote the time-space cylinder by  $Q = \Omega \times (0, T)$ , its boundary by  $\Sigma = \Gamma \times (0, T)$ .
- (ii) The initial value  $y_0$  is a given function in  $H_g$ . The desired state has to satisfy  $y_T \in H_g$ .
- (iii) The coefficient  $\alpha_T$  and the regularization parameter  $\gamma$  are positive real numbers.
- (iv) The control constraints  $u_a, u_b \in L^2(Q)^2$  have to satisfy  $u_{a,i}(x, t) \leq u_{b,i}(x, t)$  a.e. on  $Q$  for  $i = 1, 2$ , such that there exists admissible functions.

We define the set of admissible controls  $U_{ad}$  by

$$U_{ad} = \{u \in L^2(Q)^2 : u_{a,i}(x, t) \leq u_i(x, t) \leq u_{b,i}(x, t) \text{ a.e. on } Q, i = 1, 2\}.$$

Then  $U_{ad}$  is non-empty, convex and closed in  $L^2(Q)^2$ .

So we end up with the optimization problem in function space

$$\min J(y, u)$$

subject to the state equation

$$\begin{aligned} y_t + \nu Ay + \nu Cy + B(y) &= u & \text{in } L^2(0, T; V'_g), \\ y(0) &= y_0 \end{aligned} \quad (3.3)$$

and the control constraint

$$u \in U_{ad}. \quad (3.4)$$

### 3.2. Existence of optimal solutions

We call a couple  $(y, u)$  of state and control *admissible* if it satisfies the constraints (3.3) - (3.4) of the optimal control problem. We will denote in the sequel pairs of control and state by  $v$ , e.g.  $v = (y, u)$ ,  $\bar{v} = (\bar{y}, \bar{u})$ , and so on.

First, we will show the existence of optimal solutions.

**Theorem 3.1.** *The optimal control problem admits a globally optimal solution  $\bar{u} \in U_{ad}$  with associated state  $\bar{y} \in W(0, T)$ .*

*Proof.* The set of admissible controls is non-empty and bounded in  $L^2(Q)^2$ . For every control in  $L^2(Q)^2$ , there exists a unique solution of the state equation (3.3), see Theorem 2.1. Furthermore, the functional  $J$  is bounded from below,  $J(y, u) \geq 0$ . Hence, there exists the infimum of  $J$  over all admissible controls and states

$$0 \leq \bar{J} = \inf_{(y,u)} J(y, u) \leq \infty.$$

Moreover, there is a minimizing sequence  $(y_n, u_n)$  of admissible pairs such that  $J(y_n, u_n) \rightarrow \bar{J}$  for  $n \rightarrow \infty$ . The set  $\{u_n\}$  is bounded in  $L^2(Q)^2$ , then we can extract a subsequence  $(y_{n'}, u_{n'})$  converging weakly in  $W(0, T) \times L^2(Q)^2$  to some limit  $(\bar{y}, \bar{u})$ . The set of admissible controls is convex and closed in  $L^2(Q)^2$ , hence it is weakly closed in  $L^2(Q)^2$ . Thus, the control  $\bar{u}$  is admissible,  $\bar{u} \in U_{ad}$ .

Now, we have to show that the pair  $(\bar{y}, \bar{u})$  satisfies the state equation (3.3). We find for  $v \in L^2(0, T; V_g)$  the convergences

$$\begin{aligned} \langle y_{n', t}, v \rangle_{L^2(V'_g), L^2(V_g)} &\rightarrow \langle \bar{y}_t, v \rangle_{L^2(V'_g), L^2(V_g)}, \\ \langle Ay_{n'}, v \rangle_{L^2(V'_g), L^2(V_g)} &\rightarrow \langle A\bar{y}, v \rangle_{L^2(V'_g), L^2(V_g)}, \\ \langle Cy_{n'}, v \rangle_{L^2(V'_g), L^2(V_g)} &\rightarrow \langle C\bar{y}, v \rangle_{L^2(V'_g), L^2(V_g)}, \\ \langle u_{n'}, v \rangle_{L^2(V'_g), L^2(V_g)} &\rightarrow \langle \bar{u}, v \rangle_{L^2(V'_g), L^2(V_g)}, \quad \text{as } n' \rightarrow \infty. \end{aligned}$$

The mapping  $w \mapsto w(0)$  is linear and continuous from  $W(0, T) \rightarrow H_g$ . Thus,  $y_n(0)$  converges weakly to  $\bar{y}(0)$ . We have  $y_0 = y_n(0)$  for all  $n$ , hence it holds  $\bar{y}(0) = y_0$ . On the other hand, we have  $y_{n'} \rightarrow \bar{y}$  strong in  $L^2(Q)^2$ , because the space  $W(0, T)$  is compactly imbedded into  $L^2(Q)^2$ , then by using [7, Lemma 3.2], we obtain

$$\langle B(y_{n'}), v \rangle_{L^2(V'_g), L^2(V_g)} \rightarrow \langle B(\bar{y}), v \rangle_{L^2(V'_g), L^2(V_g)}, \quad \text{for } n' \rightarrow \infty.$$

Consequently, all addends in the weak formulation of the state equation converge, and

$$\langle \bar{y}_t + \nu A\bar{y} + \nu C\bar{y} + B(\bar{y}) - \bar{u}, v \rangle_{L^2(V'_g), L^2(V_g)} = 0$$

is fulfilled for all  $v \in L^2(0, T; V_g)$ . Furthermore, the initial condition  $\bar{y}(0) = y_0$  is satisfied. Hence,  $\bar{y}$  is itself the weak solution of the state equation with right-hand side  $\bar{u}$ , i.e.  $\bar{y} = S(\bar{u})$ .

Finally, it remains to show  $\bar{J} = J(\bar{y}, \bar{u})$ . The objective functional consists of several norm squares, thus it is weakly lower semicontinuous which implies

$$J(\bar{y}, \bar{u}) \leq \liminf J(y'_{n'}, u'_{n'}) = \bar{J}, \quad \text{as } n' \rightarrow \infty.$$

Since  $(\bar{y}, \bar{u})$  is admissible, and  $\bar{J}$  is the infimum over all admissible pairs, it follows that  $\bar{J} = J(\bar{y}, \bar{u})$ . The proof is complete.  $\square$

### 3.3. Lagrange functional

We define the Lagrange functional  $\mathcal{L} : W(0, T) \times L^2(Q)^2 \times L^2(0, T; V_g) \rightarrow \mathbb{R}$  for the optimal control problem as follows:

$$\mathcal{L}(y, u, \lambda) = J(y, u) - \langle y_t + \nu Ay + \nu Cy + B(y) - u, \lambda \rangle_{L^2(V'_g), L^2(V_g)}.$$

The Lagrange function  $\mathcal{L}$  is, for given  $\lambda \in L^2(0, T; V_g)$ , twice Fréchet-differentiable with respect to  $(y, u) \in W(0, T) \times L^2(Q)^2$ . The first-order derivatives of  $\mathcal{L}$  with respect to  $y$  and  $u$  in direction  $w \in W(0, T)$  and  $h \in L^2(Q)^2$  respectively are

$$\begin{aligned} \mathcal{L}_y(y, u, \lambda)w &= \alpha_T(y(T) - y_T, w(T))_{H_g} \\ &\quad - \langle w_t + \nu Aw + \nu Cw + B'(y)w, \lambda \rangle_{L^2(V'_g), L^2(V_g)}. \end{aligned} \quad (3.5)$$

$$\mathcal{L}_u(y, u, \lambda)h = \gamma(u, h)_Q + (h, \lambda)_Q. \quad (3.6)$$

The second-order derivative of  $\mathcal{L}$  with respect to  $v = (y, u)$  in directions  $(w_1, h_1), (w_2, h_2) \in W(0, T) \times L^2(Q)^2$  is

$$\mathcal{L}_{vv}(y, u, \lambda)[(w_1, h_1), (w_2, h_2)] = \mathcal{L}_{yy}(y, u, \lambda)[w_1, w_2] + \mathcal{L}_{uu}(y, u, \lambda)[h_1, h_2],$$

with

$$\mathcal{L}_{yy}(y, u, \lambda)[w_1, w_2] = \alpha_T(w_1(T), w_2(T))_{H_g} - \langle B''(y)[w_1, w_2], \lambda \rangle_{L^2(V_g), L^2(V_g)},$$

and

$$\mathcal{L}_{uu}(y, u, \lambda)[h_1, h_2] = \gamma(h_1, h_2)_Q.$$

**Theorem 3.2.** [11] Let  $\lambda \in L^2(0, T; V_g)$  be given. Then the Lagrangian  $\mathcal{L}$  is twice Fréchet-differentiable with respect to  $v = (y, u)$  from  $W(0, T) \times L^2(Q)^2$  to  $\mathbb{R}$ . The second-order derivative at  $(y, u)$  fulfills, together with the Lagrange multiplier  $\lambda$ , the estimate

$$|\mathcal{L}_{yy}(y, u, \lambda)[w_1, w_2]| \leq c_{\mathcal{L}}(1 + \|\lambda\|)\|w_1\|\|w_2\| \quad \forall w_1, w_2 \in W(0, T),$$

with some constant  $c_{\mathcal{L}} > 0$  that does not depend on  $y, u, \lambda, w_1, w_2$ .

**Remark 3.1.** The objective functional  $J$  is twice continuously differentiable from  $W(0, T) \times L^2(Q)^2$  to  $\mathbb{R}$ . The reduced objective  $\phi(u) := J(S(u), u)$  continuously is also twice continuous differentiable from  $L^2(Q)^2$  to  $\mathbb{R}$ .

#### 4. First-order necessary optimality condition

**Definition 4.1.** [11] A control  $\bar{u} \in U_{ad}$  is said to be locally optimal in  $L^p(Q)^2$ , if there exists a constant  $\rho > 0$  such that

$$J(\bar{y}, \bar{u}) \leq J(y, u)$$

holds for all  $u \in U_{ad}$  with  $\|\bar{u} - u\|_p \leq \rho$ . Here,  $\bar{y}$  and  $y$  denote the states associated with  $\bar{u}$  and  $u$ , e.g.  $\bar{y} = S(\bar{u})$  and  $y = S(u)$ .

Now, we will state and prove the first-order necessary optimality condition.

**Theorem 4.1.** Let  $\bar{u}$  be locally optimal in  $L^2(Q)^2$  with associated state  $\bar{y} = S(\bar{u})$ . Then there exists a unique Lagrange multiplier  $\bar{\lambda} \in W^{4/3}(0, T; V_g)$ , which is the weak solution of the adjoint equation

$$\begin{aligned} -\bar{\lambda}_t + \nu A\bar{\lambda} + \nu C\bar{\lambda} + B'(\bar{y})^* \bar{\lambda} &= 0, \\ \bar{\lambda}(T) &= \alpha_T(\bar{y}(T) - y_T). \end{aligned} \quad (4.1)$$

Moreover, the variational inequality

$$(\gamma\bar{u} + \bar{\lambda}, u - \bar{u})_Q \geq 0 \quad \forall u \in U_{ad} \quad (4.2)$$

is satisfied.

*Proof.* From Remark 3.1,  $\phi$  is twice Fréchet-differentiable from  $L^2(Q)^2$  to  $\mathbb{R}$ . By Banach space optimization principles, we know that the variational inequality

$$\phi'(\bar{u})(u - \bar{u}) \geq 0 \quad \forall u \in U_{ad} \quad (4.3)$$

is necessary for local optimality of  $\bar{u}$ . It remains to compute  $\phi'$  and to derive the adjoint system. Let us write  $\phi$  in the following way

$$\phi(u) = \frac{\alpha_T}{2} |FS(u) - y_T|_2^2 + \frac{\gamma}{2} \|u\|_2^2,$$

where  $F$  defined by  $F(y) = y(T)$  is a linear and continuous operator from  $W(0, T)$  to  $H_g$ . The first derivative  $\phi'$  at  $\bar{u}$  in direction  $h \in L^2(Q)^2$  is characterized by

$$\phi'(\bar{u})h = \alpha_T(FS(\bar{u}) - y_T, FS'(\bar{u})h)_2 + \gamma(\bar{u}, h)_Q.$$

Now, we compute the adjoint operator of  $F$ .

First, we have by definition  $\langle F^*v, w \rangle = (v, w(T))_2$  for  $w \in W(0, T)$  and  $v \in H_g$ . This implies with  $v := Fy$  that  $\langle F^*Fy, w \rangle = (y(T), w(T))_2$  holds for all  $y, w \in W(0, T)$ . With the help of this adjoint, we transform  $\phi'(\bar{u})h$  to

$$\phi'(\bar{u})h = \langle \alpha_T F^*(FS(\bar{u}) - y_T), S'(\bar{u})h \rangle + \gamma(u, h)_Q.$$

Substituting back  $S(\bar{u}) = \bar{y}$ , the expression

$$\phi'(\bar{u})h = \langle \alpha_T F^*(\bar{y}(T) - y_T), S'(\bar{u})h \rangle + \gamma(u, h)_Q. \quad (4.4)$$

is found. The state  $\bar{y}$  is solution of the nonlinear state equation, therefore it belongs to  $W(0, T)$ . Altogether, the addends defining  $\phi'(\bar{u})$  have the following regularity:

$$F^*(\bar{y}(T) - y_T) \in W(0, T)^*.$$

Let us define

$$\bar{\lambda} := S'(\bar{u})^* \{ \alpha_T F^*(\bar{y}(T) - y_T) \}. \quad (4.5)$$

The term in brackets fits in the assumption of Lemma 2.4. Thus, we can interpret  $\bar{\lambda}$  as the weak solution of

$$\begin{aligned} -\lambda_t + \nu A(\lambda) + \nu C(\lambda) + B'(\bar{y})^* \lambda &= 0 \quad \text{in } L^{4/3}(0, T; V'_g), \\ \lambda(T) &= \alpha_T (\bar{y}(T) - y_T). \end{aligned}$$

Furthermore, we get the regularity  $\lambda \in W^{4/3}(0, T)$ . The space  $W^{4/3}(0, T)$  is continuously imbedded into  $L^p(Q)^2$  for  $p < 7/2$ , see [10]. Combining the definition of  $\bar{\lambda}$  in (4.5) with (4.3) and (4.4) yields

$$0 \leq \phi'(\bar{u})(u - \bar{u}) = (\gamma \bar{u} + \bar{\lambda}, u - \bar{u})_Q \quad \forall u \in U_{ad},$$

which is the variational inequality (4.2). □

The necessary optimality conditions can be formulated equivalently using the Lagrangian.

**Theorem 4.2.** *Under the conditions of Theorem 4.1, it is necessary for local optimality of  $\bar{u}$  that there exists  $\bar{\lambda} \in W^{4/3}(0, T)$  such that*

$$\begin{aligned}\mathcal{L}_y(\bar{y}, \bar{u}, \bar{\lambda})w &= 0 \quad \forall w \in W_0, \\ \mathcal{L}_y(\bar{y}, \bar{u}, \bar{\lambda})(u - \bar{u}) &\geq 0 \quad \forall u \in U_{ad}.\end{aligned}$$

are fulfilled.

*Proof.* Let  $\bar{\lambda}$  be solution of (4.1), which is given by (4.5). According to (3.5), we can write  $\mathcal{L}_y(\bar{y}, \bar{u}, \bar{\lambda})w$  as

$$\mathcal{L}_y(\bar{y}, \bar{u}, \bar{\lambda})w = \alpha_T(\bar{y}(T) - y_T, w(T))_{H_g} - \langle w_t + \nu Aw + \nu Cw + B'(y)w, \bar{\lambda} \rangle_{L^2(V'_g), L^2(V_g)}.$$

By construction of  $\bar{\lambda}$ , the right-hand side vanishes for  $w \in W_0$ , as Lemma 2.3. Hence, we obtain  $\mathcal{L}_y(\bar{y}, \bar{u}, \bar{\lambda})w = 0$  for all test functions  $w \in W_0$ . From (3.6), the derivative of the Lagrangian with respect to the control  $u$  is given as

$$\mathcal{L}_y(\bar{y}, \bar{u}, \bar{\lambda})(u - \bar{u}) = (\gamma \bar{u} + \bar{\lambda}, u - \bar{u})_Q.$$

Consequently, this expression must be greater or equal to zero for all admissible control  $u \in U_{ad}$ .  $\square$

**Remark 4.1.** One may use the approach as in [9], [12] to study the second-order optimality conditions as well as the regularity and stability of optimal controls for 2D  $g$ -Navier-Stokes equations. This is the subject of our future work.

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*Manuscript received 25-07-2016; accepted 14-12-2016.*



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